

Block 16 - Ampère's Law and Magnetomotive Force (MMF)

Student Group

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16.0 Intro

16.0.1 Learning objectives

After this 90-minute block, you can

- ...

16.0.2 Preparation at Home

Well, again

- read through the present chapter and write down anything you did not understand.
- Also here, there are some clips for more clarification under 'Embedded resources' (check the text above/below, sometimes only part of the clip is interesting).

For checking your understanding please do the following exercises:

- ...

16.0.3 90-minute plan

1. Warm-up (x min):
 1.
2. Core concepts & derivations (x min):
 1. ...
3. Practice (x min): ...
4. Wrap-up (x min): Summary box; common pitfalls checklist.

16.0.4 Conceptual overview

1. ...

16.1 Core content

16.1.1 Generalization of the Magnetic Field Strength

So far, only the rotational symmetric problem of a single wire was considered in formula. I.e a current I and the length s of a magnetic field line around the wire was given to calculate the magnetic field strength H :

$$\begin{aligned} \quad H_{\varphi} = \frac{I}{2 \cdot \pi \cdot r} \quad \Leftrightarrow \\ \quad I = H_{\varphi} \cdot s \quad \quad \quad | \quad \text{\textit{applies only to the long, straight conductor}} \end{aligned}$$

Now, this shall be generalized. For this purpose, we will look back at the electric field.

For the electric field strength E of a capacitor with two plates at a distance of s and the potential difference U holds:

$$\begin{aligned} U = E \cdot s \quad \quad | \quad \text{\textit{applies to plate capacitor only}} \end{aligned}$$

In words: The potential difference is given by adding up the field strength along the path of a probe from one plate to the other.

This was extended to the voltage between to points 1 and 2 . Additionally, we know by Kirchhoff's voltage law that the voltage on a closed path is "0".

$$\begin{aligned} U_{12} &= \int_1^2 \vec{E} \cdot d\vec{s} \quad \quad U = \oint \vec{E} \cdot d\vec{s} \\ &= 0 \end{aligned}$$

We can now try to look for similarities. Also for the magnetic field, the magnitude of the field strength is summed up along a path to arrive at another field-describing quantity.

Because of the similarity the so-called **magnetic potential difference V_m** between point 1 and 2 is introduced:

$$\begin{aligned} V_m &= H \cdot s \quad \quad | \quad \text{\textit{applies to rotational symmetric problems only}} \end{aligned}$$

$$\boxed{V_m = V_{m, 12} = \int_1^2 \vec{H} \cdot d\vec{s} \quad \quad V_m = \oint \vec{H} \cdot d\vec{s} = \theta}$$

We need to take a closer look here. Any closed path in the static electric field leads to a potential difference of $U = \oint \vec{E} \cdot d\vec{s} = 0$.

BUT: closed paths in the static magnetic field leads to a magnetic potential difference which is **not mandatorily** 0 ! $V_m = \oint \vec{H} \cdot d\vec{s} = \theta$

Another new quantity is introduced: the **magnetic voltage θ** :

1. The magnetic voltage θ is the magnetic potential difference on a closed path.
2. Since the magnetic voltage θ is valid for exactly one turn along our single wire, θ is also equal to the current through the wire:

$$\begin{aligned} \theta = H \cdot s = I \quad \quad | \quad \text{\textit{applies only to the long, straight conductor}} \end{aligned}$$

3. The magnetic potential difference can take a fraction or a multiple of one turn and is therefore **not mandatorily** equal to I .
4. The magnetic voltage is generalized in the following box.

Notice:

The path integral of the magnetic field strength along an arbitrary closed path is equal to the free currents (= current density) through the surface enclosed by the path.

The magnetic voltage θ (and therefore the current) is the cause of the magnetic field strength.

This leads to the **Ampere's Circuital Law**

$\boxed{\oint_{\mathcal{S}} \vec{H} \cdot d\vec{s} = \theta}$	<p>The magnetic voltage θ can be given as</p> <ul style="list-style-type: none"> • $\theta = I$ for a single conductor • $\theta = N \cdot I$ for a coil • $\theta = \sum_n I_n$ for multiple conductors • $\theta = \int_A \vec{S} \cdot d\vec{A}$ for any spatial distribution (see block15)
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The unit of the magnetic voltage θ is **Ampere** (or **Ampere-turns**).

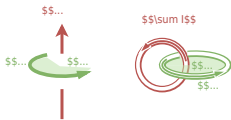
In the english literature the magnetic voltage is called **Magnetomotive force**

Notice:

$\oint \vec{s}$ and $\oint \vec{A}$ in $\oint \vec{H} \cdot d\vec{s} = \theta = \int_A \vec{S} \cdot d\vec{A}$ build a right-hand system.

1. Once the thumb of the right hand is pointing along $\oint \vec{A}$, the fingers of the right hand show the correct direction for $\oint \vec{s}$ for positive \vec{H} and \vec{S}
2. Currents into the direction of the right hand's thumb count positive. Currents antiparallel to it count negative.

Fig. 1: Right hand rule



16.1.2 Recap of the fieldline images

longitudinal coil

Fig. 2: Magnetic field in a longitudinal coil

A longitudinal coil can be seen in [figure 2](#).

The created field density of the coil can be derived from Ampere's Circuital Law

$$\begin{aligned} \oint \vec{H}(t) \cdot d\vec{s} &= \int_{\text{inner}} \vec{H}_{\text{inner}}(t) \cdot d\vec{s} + \int \vec{H}_{\text{outer}}(t) \cdot d\vec{s} \\ &= \int \vec{H}(t) \cdot d\vec{s} + 0 \\ &= H(t) \cdot l \end{aligned}$$

The magnetic field in a toroidal coil is often considered as homogenous in the inner volume, when the length l is much larger than the diameter: $l \gg d$.

With a given number N of windings, the magnetic field strength H is

toroidal coil

Fig. 3: Magnetic field in a toroidal coil

A toroidal coil has a donut-like setup. This can be seen in [figure 3](#).

The toroidal coil is often defined by:

- The minor radius r : The radius of the circular cross-section of the coil.
- The major radius R : The distance from the center of the entire toroid (the center of the hole) to the center of the circular cross-section of the coil.

For reasons of symmetry, it shall get clear that the field lines form concentric circles. Also the magnetic field strength H in a toroidal coil is often considered as homogenous, when $R \gg r$. With a given number N of windings, the magnetic field strength H is

$$\begin{aligned} \theta &= H \cdot l = N \cdot I \\ \boxed{H} &= \frac{N \cdot I}{l} \quad \text{longitudinal coil} \end{aligned}$$

$$\begin{aligned} \theta &= H \cdot 2\pi R = N \cdot I \\ \boxed{H} &= \frac{N \cdot I}{2\pi R} \quad \text{toroidal coil} \end{aligned}$$

16.2 Common pitfalls

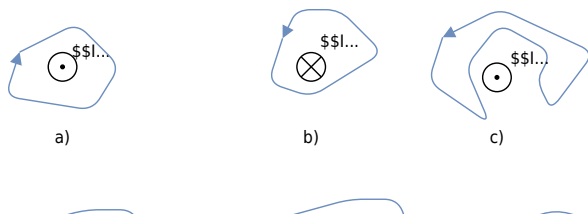
- ...

16.3 Exercises

Task 3.2.3 Magnetic Potential Difference

Fig. 3: different trajectories around current-carrying conductors
 Result e)

e) $V_{(r, \theta)} = 20 \ln \left(\frac{2.25}{2.5} \right) \text{ Am} = -2.5 \text{ Am}$



Given are the adjacent closed trajectories in the magnetic field of current-carrying conductors (see figure 3). Let $I_1 = 2 \text{ A}$ and $I_2 = 4.5 \text{ A}$ be valid.

In each case, the magnetic potential difference V_{m} along the drawn path is sought.

Path

- The magnetic potential difference is given as the **sum of the current through the area within a closed path**.
- The direction of the current and the path have to be considered with the righthand rule.

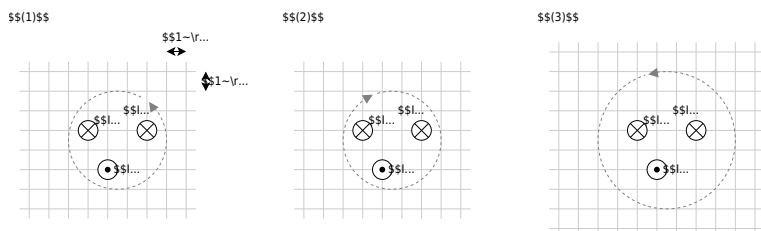
Exercise E7 Magnetic Voltage (written test, approx. 6 % of a 120-minute written test, SS2021)

The following images show cross-sections of electrical cables.

A closed path is shown as a dashed line. The magnetic voltage θ on these paths shall be analyzed.

The following values are given for the currents:

- $\begin{aligned} \theta_{(1)} &= -4 \text{ A} \\ \theta_{(2)} &= 0 \text{ A} \\ \theta_{(3)} &= 5 \text{ A} \\ \theta_{(4)} &= 5 \text{ A} \end{aligned}$
- $I_3 = 1 \text{ A}$
 - $I_4 = 4 \text{ A}$



Specify which magnetic voltages $\theta_{(1)}$, $\theta_{(2)}$, and $\theta_{(3)}$ result. Note the direction of the path in each case!

Path

For the resulting current the direction of the path has to be considered with the right-hand rule:

- $$I_{(1)} = +I_2 - I_1 - I_3 \quad \rightarrow \quad \theta_{(1)} = 2 \text{ A} - 5 \text{ A} - 1 \text{ A}$$
- $$I_{(2)} = +I_3 + I_4 - I_1 \quad \rightarrow \quad \theta_{(2)} = 1 \text{ A} + 4 \text{ A} - 5 \text{ A}$$
- $$I_{(3)} = +I_3 - I_4 - I_2 \quad \rightarrow \quad \theta_{(3)} = 1 \text{ A} - 4 \text{ A} - 2 \text{ A}$$

Exercise E10 Magnetic Potential

(written test, approx. 8 % of a 120-minute written test, SS2024)

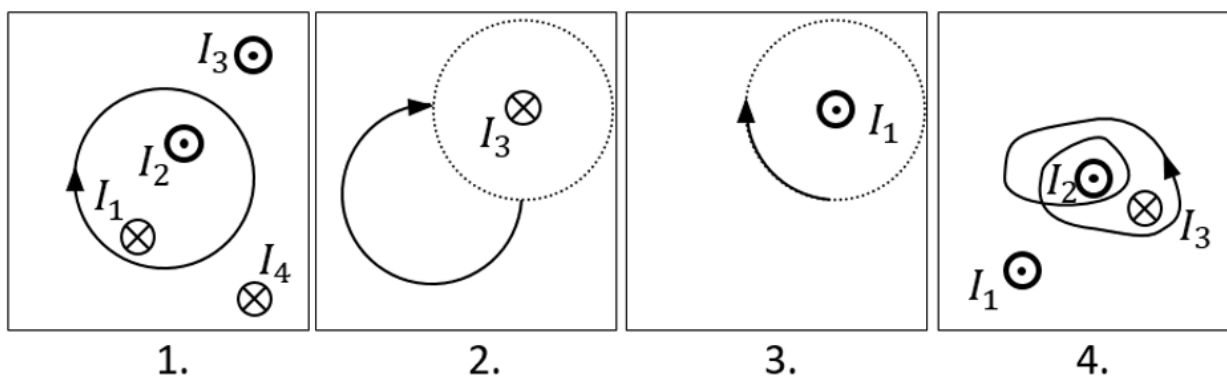
Calculate the magnetic potential difference V_{m} for the following paths as shown by the solid lines.

Dotted lines are only for there for symmetry aspects!

The wires conduct the following currents:

- $|I_1| = 2 \text{ A}$
- $|I_2| = 5 \text{ A}$
- $|I_3| = 11 \text{ A}$
- $|I_4| = 7 \text{ A}$

Pay attention to the signs of the currents (given by the diagrams) and of the results!



Result

Based on the right-hand rule and the part of a full revolution the following results:

1. Task: $+I_1 - I_2 = -3 \text{ A}$
2. Task: $+\frac{1}{4} I_3 = 11/4 \text{ A}$ (it does not matter which way the path goes from the startpoint to the endpoint, as long as it has the same direction and number of revolutions)
3. Task: $-\frac{1}{4} I_1 = -0.5 \text{ A}$
4. Task: $+2 \cdot I_2 - 1 \cdot I_3 = -1 \text{ A}$

Exercise E6 Fields of an coax Cable
(written test, approx. 12 % of a 120-minute written test, SS2024)

2. With the graph of the magnitude of D (in digital bar fields) DC (the dia $\phi = 0.6 \text{ mm}$) shows the cross-section of the coax cable with (0,6 origin | 0) since the diagram is a top-down view and label for the diagram appears:

Path

- Inner conductor: $+3.3 \text{ mA}$, $+10 \text{ nC}$ (current into the plane of the path diagram)
- Outer conductor: -3.3 mA , 0 nC (current out of the plane of diagram)
- for $(0.1 \text{ mm} | 0)$: $B_{\text{I}} = 5.28 \dots \text{ A/m}^2$
- for $(0.55 \text{ mm} | 0)$: $B_{\text{O}} = 6.985 \dots \text{ A/m}^2$

The magnitude of the electric displacement field D can be calculated by: $\int D \cdot dA = Q_{\text{enc}}$.

Here, in any position r from the center, the surrounding area is the surface of a cylindrical shape (here for simplicity without the round endings).

Since the charge is spread on the surface of the conductor, it is not so much about the volume, but about the surface. This leads to: $D(x) = \frac{Q}{A} = \frac{Q}{2\pi r \cdot \Delta z}$

This is proportional to the area within this radius. Therefore, the formula $H = \frac{I}{2\pi r}$ is also valid for D . So, we get for $D_{\text{I}}(r)$ at $r = 0.1 \text{ mm}$, and $D_{\text{O}}(r)$ at $r = 0.55 \text{ mm}$.

For x within the outer conductor one also gets a linear proportionality with a similar approach: $D = \frac{Q}{2\pi r \cdot \Delta z} = \frac{10 \cdot 10^{-9} \text{ C}}{2\pi \cdot \{0.1 \cdot 10^{-3} \text{ m}\} \cdot 0.5 \text{ m}} \parallel D_{\text{I}}(r) = \frac{Q}{2\pi r \cdot \Delta z} = \frac{10 \cdot 10^{-9} \text{ C}}{2\pi \cdot \{0.55 \cdot 10^{-3} \text{ m}\} \cdot 0.5 \text{ m}} \parallel$

Hint: For the direction, one has to consider the sign of the enclosed charge. By this, we see that the D -field is positive.

But here, again only the magnitude was questioned!

.. What is the magnitude of the magnetic field strength H at $(0.1 \text{ mm} | 0)$ and $(0.55 \text{ mm} | 0)$?

Path

The magnitude of the magnetic field strength H can be calculated by: $H = \frac{I}{2\pi r}$

So, we get for H_{I} at $(0.1 \text{ mm} | 0)$, and H_{O} at $(0.55 \text{ mm} | 0)$

~mm | 0)\$:

$$\begin{aligned} H_{\text{i}} &= \frac{I}{2 \pi \cdot r_{\text{i}}} \quad \&= \frac{+3.3 \text{ A}}{2 \pi \cdot \{0.1 \cdot 10^{-3} \text{ m}\}} \\ H_{\text{o}} &= \frac{I}{2 \pi \cdot r_{\text{o}}} \quad \&= \frac{+3.3 \text{ A}}{2 \pi \cdot \{0.55 \cdot 10^{-3} \text{ m}\}} \end{aligned}$$

Hint: For the direction, one has to consider the right-hand rule. By this, we see that the H -field on the right side points downwards.

Therefore, the sign of the H -field is negative.

But here, only the magnitude was questioned!

Embedded resources

Explanation (video): ...

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Permanent link:

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